

# Laser Magnetic Resonance Spectroscopy of ClO and Kinetic Studies of the Reactions of ClO with NO and NO<sub>2</sub>

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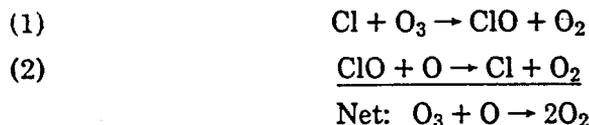
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## Abstract

Far-infrared rotational transitions in ClO( $X^2\Pi_{3/2}, v = 0$ ) have been observed using laser magnetic resonance (LMR) with an optically pumped spectrometer. Five observed transitions at wavelengths between 444 and 713  $\mu\text{m}$  have been compared with values predicted with spectroscopic constants from the literature. LMR detection of ClO has been used to study its reactions with NO and NO<sub>2</sub> in a discharge flow system under pseudo-first-order conditions for ClO. The measured rate constants are  $k(\text{ClO} + \text{NO}) = (7.1 \pm 1.4) \times 10^{-12} \exp\{(270 \pm 50)/T\} \text{ cm}^3/\text{molec}\cdot\text{s}$  for the temperature range of  $202 < T < 393 \text{ K}$ ;  $k(\text{ClO} + \text{NO}_2 + \text{M}) = (2.8 \pm 0.6) \times 10^{-33} \exp\{(1090 \pm 80)/T\} \text{ cm}^6/\text{molec}^2\cdot\text{s}$  ( $\text{M} = \text{He}, 250 < T < 387 \text{ K}$ ),  $(3.5 \pm 0.6) \times 10^{-33} \exp\{(1180 \pm 80)/T\}$  ( $\text{M} = \text{O}_2, 250 < T < 416 \text{ K}$ ), and  $(2.09 \pm 0.3) \times 10^{-31}$  ( $\text{M} = \text{N}_2, T = 297 \text{ K}$ ). All measurements were made at low pressures, between 0.6 and 6.6 torr. These results are compared with those from other studies.

## Introduction

The potential impact of stratospheric ozone depletion by the chlorine catalytic cycle



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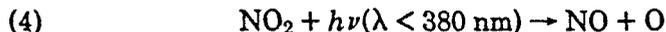
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has generated considerable interest [1]. To evaluate the effects of chlorine on ozone one must consider the reactions involving the active species Cl and ClO. The presence of nitrogen oxides (NO and NO<sub>2</sub>) in the stratosphere limit the effectiveness of the chlorine cycle. First, the reaction with NO,



when combined with



and reaction (1) produces no net destruction of ozone. Second, the reaction of ClO with NO<sub>2</sub>,

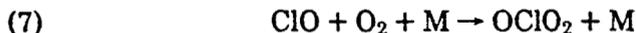


forms ClONO<sub>2</sub> [2], which is thought to be an unreactive reservoir for chlorine.

There have been five studies of the ClO + NO reaction. Clyne and Watson [3] first measured the rate constant  $k = (1.7 \pm 0.2) \times 10^{-11} \text{ cm}^3/\text{molec}\cdot\text{s}$  at room temperature using discharge-flow/mass spectrometry in 1974. Later Zahniser and Kaufman [4] employed the discharge-flow/resonance fluorescence technique to determine the rate constant relative to that of the Cl + O<sub>3</sub> reaction. They reported rate constants at temperatures ranging from 230 to 295 K, giving a room temperature value slightly higher than the previous value and a negative Arrhenius activation energy of 400 cal/mol (1 cal = 4.184 J). Leu and DeMore [5] investigated the temperature dependence (227–415 K) using mass spectrometry and obtained slightly lower rate constants and a negative activation energy of about 590 cal/mol. Recently Clyne and MacRobert [6] and Ray and Watson [7] repeated the measurement at room temperature and reported rate constants close to the first study.

The ClO + NO<sub>2</sub> + M reaction in the low-pressure range of 1–6 torr (1 torr = 133.3 Pa) has been studied by three groups [8–10] using discharge-flow techniques with various detection methods. Leu et al. [8] and Birks et al. [9] detected ClO by mass spectrometry. This technique was complicated by effects of ClONO<sub>2</sub> fragmentation in the mass spectrometer ion source to produce ClO. Since the rate constant measurement is based on the pseudo-first-order decay of ClO, this effect could lead to an underestimation of the true value of  $k_6$ . Birks et al. used an empirically derived correction for this effect that increased the measured value by up to about 17%, while Leu et al. operated under conditions which they stated required no correction. Zahniser et al. [10] detected ClO indirectly by conversion of ClO to Cl with NO addition, reaction (3), followed by resonance fluorescence detection of the Cl atoms.

While the various rate constant measurements for both reactions (3) and (6) are in reasonable agreement, some significant discrepancies still exist. For the  $\text{ClO} + \text{NO}$  reaction, the rate constants given by the two temperature dependence studies [4,5] differ by about 40%. For the  $\text{ClO} + \text{NO}_2 + \text{M}$  reaction the reported rate constants at room temperature differ by about 30% in He and by about 20% in  $\text{N}_2$ . In addition, the reaction



may take place at high  $\text{O}_2$  concentrations and could have an effect on the  $\text{ClO} + \text{NO}_2 + \text{O}_2$  reaction. Since no measurements of  $k_6$  were made in  $\text{O}_2$ , it is interesting to investigate reaction (6) in  $\text{O}_2$  as a first step toward the understanding of the importance of reaction (7) and the chemistry of  $\text{OClO}_2$  in the stratosphere. Also it is worthwhile to test a new method which detects  $\text{ClO}$  directly, has high sensitivity, and does not have possible complications from the product  $\text{ClNO}_3$  or other species.

Here we report the far-infrared (FIR) laser magnetic resonance (LMR) spectra of  $\text{ClO}$  at five different laser frequencies. Applying this LMR detection technique, we have also investigated the temperature dependence of the rate constants of reactions (3) and (6).

### Experimental

An LMR spectrometer was used in conjunction with a discharge-flow tube in this study. The basic principles of LMR for spectroscopic as well as kinetic studies have been described previously [11,12]. Briefly, a continuous-wave  $\text{CO}_2$  laser with an output power of approximately 25 W was used to pump a FIR  $\text{CH}_3\text{OH}$ ,  $\text{CD}_3\text{I}$ , or  $\text{DCOOH}$  laser which produces about 100  $\mu\text{W}$ . The intracavity absorption was accomplished by Zeeman tuning a rotational transition of the  $\text{ClO}$  molecule into resonance with the FIR laser frequency. The magnetic field can be varied up to 21 kG. High sensitivity was obtained by ac modulating the magnetic field and using frequency-locked phase-sensitive detection. It has been found that the peak-to-peak height of the first derivative of the absorption curve is proportional to the radical concentration [11].

In the kinetic studies, the rate constants were measured by observing the pseudo-first-order decay of  $\text{ClO}$  concentrations in a large excess of  $\text{NO}$  or  $\text{NO}_2$  as a function of reaction time. The  $\text{ClO}$  radical was detected at  $\approx 0.5$  kG using a  $\pi$ -polarized 556.9- $\mu\text{m}$   $\text{CD}_3\text{I}$  laser line pumped by the 10P(36)  $\text{CO}_2$  laser emission. The detection limit for  $\text{ClO}$  is approximately  $2 \times 10^8$  molec/ $\text{cm}^3$  (signal-to-noise ratio  $\approx 1$  at 1 s time constant).

The discharge-flow apparatus [12] consists of a 2.54-cm inside diameter, 110-cm long Pyrex flow tube with a 9-mm outside diameter movable inlet. The flow tube was mounted in the LMR spectrometer with the tube exit positioned directly above the intersection of the FIR laser beam and the magnetic field.

The total gas flow consisted of the carrier gas (He, O<sub>2</sub>, or N<sub>2</sub>), NO<sub>2</sub> or NO, and the gas mixture from the ClO source. For the ClO + NO<sub>2</sub> + M reaction rate measurements, the ClO source was part of the movable inlet which allowed variation of the position of the ClO injection into the flow tube. The ClO radical was formed by reacting Cl atoms with excess of O<sub>3</sub> in the top part of the movable inlet



The Cl was generated by passing a trace of Cl<sub>2</sub>  $\approx$  0.002 STP cm<sup>3</sup>/s (STP = 273 K, 1 atm) in helium carrier gas (0.2–0.4 STP cm<sup>3</sup>/s) through a 20–35-W microwave discharge. Ozone was eluted with helium from a 196 K silica gel trap and diluted with additional helium (0.3–1.0 STP cm<sup>3</sup>/s) before reacting with Cl atoms. The ClO formation is >99% completed within the first 5-cm length of injector based on the reported value [13] of  $k_1 = 1.1 \times 10^{-11}$  cm<sup>3</sup>/molec-s and the ozone concentration of about  $3 \times 10^{14}$  molec/cm<sup>3</sup>. The distance from the O<sub>3</sub> inlet to the end of the injector was about 125 cm, long enough for any vibrationally excited ClO formed by reaction (1) to be quenched before it entered the reaction region. The ClO concentration in these experiments was about  $(1-8) \times 10^{12}$  molec/cm<sup>3</sup>. At this concentration ClO does not react significantly with itself in the flow tube [13].

For the ClO + NO<sub>2</sub> + M study, NO<sub>2</sub> was admitted into the flow tube through a side arm about 80 cm upstream of the LMR detection region. The NO<sub>2</sub> concentration in the flow tube ranged from 2 to  $20 \times 10^{14}$  molec/cm<sup>3</sup>. The NO<sub>2</sub> flow rate was measured by observing the rate of the pressure change in a calibrated volume. The effect of NO<sub>2</sub> dimerization has been taken into account in the NO<sub>2</sub> flow rate measurements by applying the correction  $F_{\text{NO}_2} = F'_{\text{NO}_2}(1 + 2\bar{P}/K_p)$ , where  $F'_{\text{NO}_2}$  is the measured apparent NO<sub>2</sub> flow rate,  $\bar{P}$  (in torr) is the average pressure in the calibrated volume, and  $K_p = [\text{NO}_2]^2/[\text{N}_2\text{O}_4]$  is the equilibrium constant (in torr) for the NO<sub>2</sub> dimerization reaction. At 298 K,  $K_p = 106$  torr [14], and the NO<sub>2</sub> flow-rate correction is less than 5% with  $\bar{P} \approx 2$  torr. NO<sub>2</sub> flow measurements were always made at low pressure to minimize the correction. The effect of NO<sub>2</sub> dimerization in the flow tube was negligible (<1%) at temperatures above 270 K for  $[\text{NO}_2] < 2 \times 10^{15}$  molec/cm<sup>3</sup> ( $K_p \approx 9.7$  torr at 270 K).<sup>1</sup> However, at 250 K  $K_p$  decreased to 1.18 torr [15], and about a 4% decrease in NO<sub>2</sub> concentration due to dimerization should be taken into account for  $[\text{NO}_2] \approx 1 \times 10^{15}$  molec/cm<sup>3</sup>.

For the ClO + NO reaction, the ClO was formed in the flow tube by reacting ozone,  $(3-12) \times 10^{11}$  molec/cm<sup>3</sup>, with excess Cl atoms,  $(8-15) \times 10^{13}$  atom/cm<sup>3</sup>, which were produced by a microwave discharge of Cl<sub>2</sub> and He in a sidearm of the flow tube. NO was mixed with  $\sim 0.5$  STP cm<sup>3</sup>/s of helium and added through the movable injector. The NO concentration in the flow tube ranged from  $(3 \text{ to } 30) \times 10^{12}$  molec/cm<sup>3</sup>.

<sup>1</sup> Obtained using thermochemical data from [15].

The carrier gas was added about 85 cm upstream of the LMR detection region with a flow rate between 2 and 30 STP cm<sup>3</sup>/s, depending on the desired flow velocity and pressure. The flow rates were measured with mass flowmeters which are accurate within  $\pm 3\%$  [12]. The flow tube pressure was measured at the center of the reaction region with a capacitance manometer with an accuracy of about  $\pm 1\%$  [12]. The flow tube pressure  $P_t$  ranged from 0.6 to 6.6 torr with the average flow velocity  $\bar{v}$  ranging from 350 to 1000 cm/s for the ClO + NO<sub>2</sub> + M kinetic studies, while  $P_t = 1-3$  torr and  $\bar{v} = 1670-2730$  cm/s for the ClO + NO studies.

Temperature variability of the flow tube from 202 to 416 K was achieved by circulating either liquid N<sub>2</sub>, cooled ethanol, or heated dibutylphthalate through the flow tube jacket. The temperature of the circulating fluid was measured with a Pt resistance thermometer at the exit of the jacket, and it was shown previously that the temperature of the flow tube was uniform to  $\pm 1^\circ\text{C}$  [12].

The He (>99.9996%, analyzed) was passed through a liquid N<sub>2</sub>-cooled molecular sieve trap before use. O<sub>2</sub> and N<sub>2</sub> (>99.97%) and Cl<sub>2</sub> (>99.96%) were used without further purification. NO reactant was passed through a dry-ice-ethanol-cooled silica gel trap to remove other nitrogen oxides. NO<sub>2</sub> was synthesized by reacting purified NO (>99.0%) with excess of O<sub>2</sub>, and stored under about 2 atm of O<sub>2</sub> for at least 24 h before use. During the kinetic measurements, the pure NO<sub>2</sub> (NO < 0.01%) was kept in an ice-water bath to maintain a constant pressure of  $\sim 260$  torr.

O<sub>3</sub> was prepared by passing O<sub>2</sub> through an ac discharge and collected in a silica gel trap kept in a dry-ice-ethanol bath. The ozone partial pressure  $P_{\text{O}_3}$  was measured by passing He-O<sub>3</sub> mixtures through a 10-cm absorption cell using  $1.15 \times 10^{-17}$  cm<sup>2</sup> as the O<sub>3</sub> absorption cross section at 254 nm [16]. Thus the flow rate of O<sub>3</sub>,  $F_{\text{O}_3}$ , was calculated from the flow rate of He through the cell  $F_{\text{He}}$ , the total pressure in the cell  $P_A$ , and  $P_{\text{O}_3}$ :

$$(8) \quad F_{\text{O}_3} = F_{\text{He}} P_{\text{O}_3} / (P_A - P_{\text{O}_3})$$

The estimated random error in a rate constant measurement is about 8%. This is based on the equation  $k \propto F_T^2 T^3$  (slope) /  $P_t^3 F_R r^2$  and the estimates of 3% error in flow rate measurements  $F_T$  and  $F_R$ , 1% in temperature  $T$ , 2% in the decay plot slope, 1% in pressure  $P_t$ , and 1% in the tube radius  $r$ . However, in consideration of possible systematic errors,  $\pm 15\%$  is a more realistic estimate of the overall accuracy of the kinetic measurements.

### ClO Spectroscopy

The LMR detection of a paramagnetic radical requires a close coincidence of an FIR laser frequency with a zero-field rotational transition, because the energy range over which a molecular rotational transition can be magnetically tuned is usually less than 1 cm<sup>-1</sup> for magnetic fields of less

than 20 kG ( $10^4$  G = 1 T). In the case of ClO, the search for proper FIR laser frequencies that can yield LMR spectra is simplified because the molecular parameters for the ground state have been accurately determined [17-23]. Using the energy expression  $E_J = B(J + 1/2)^2 - D(J + 1/2)^4$  from the analysis of Kakar et al. [17] and the values of  $B = 0.61976$  cm $^{-1}$ ,  $D = 1.315 \times 10^{-6}$  cm $^{-1}$  for  $^{35}\text{ClO}$ , and  $B = 0.60930$  cm $^{-1}$ ,  $D = 1.278 \times 10^{-6}$  cm $^{-1}$  for  $^{37}\text{ClO}$ , the zero-field rotational transition energies of ClO can be estimated with sufficient accuracy to predict near resonances. Listed in Table I are the estimated energies for the  $\Delta J = +1$  rotational transitions of the ground  $^2\Pi_{3/2}$  ( $v = 0$ ) state of  $^{35}\text{ClO}$  and  $^{37}\text{ClO}$ .

The magnetic tunability of a rotational transition can be estimated from the first-order Zeeman interaction energy

$$(9) \quad E_J = -g_J M_J \mu_B H$$

where  $g_J$  is the rotational  $g$  factor,  $\mu_B$  is the Bohr magneton,  $M_J$  is the component of  $J$  in the field direction, and  $H$  is the magnetic field strength. The  $g_J$  factor for ClO may be approximated by [24]

$$(10) \quad g_J \cong \frac{(\Lambda + \Sigma)(\Lambda + 2\Sigma)}{J(J + 1)}$$

where  $J$  is the total angular momentum quantum number,  $\Lambda$  is the orbital angular momentum projection quantum number, and  $\Sigma$  is the spin angular momentum projection quantum number. Thus of the two components

TABLE I. Estimated energies and vacuum wavelengths for  $\Delta J = +1$  rotational transitions of the ground  $^2\Pi_{3/2}$  ( $v = 0$ ) state of  $^{37}\text{ClO}$  and  $^{35}\text{ClO}$ .

$J''$ (lower level)	$^{37}\text{ClO}$		$^{35}\text{ClO}$	
	$E$ (cm $^{-1}$ )	$\lambda$ ( $\mu\text{m}$ )	$E$ (cm $^{-1}$ )	$\lambda$ ( $\mu\text{m}$ )
1.5	3.046	3283	3.099	3227
2.5	4.265	2345	4.338	2305
3.5	5.483	1824	5.577	1793
4.5	6.701	1492	6.816	1467
5.5	7.919	1263	8.055	1241
6.5	9.137	1094	9.294	1076
7.5	10.35	965.7	10.53	949.4
8.5	11.57	864.1	11.77	849.5
9.5	12.79	781.9	13.01	768.7
10.5	14.01	714.0	14.25	701.9
11.5	15.22	656.9	15.48	645.8
12.5	16.44	608.3	16.72	598.1
13.5	17.65	566.4	17.96	556.9
14.5	18.87	530.0	19.19	521.0
15.5	20.08	497.9	20.43	489.5
16.5	21.30	469.5	21.66	461.6
17.5	22.51	444.2	22.90	436.7
18.5	23.72	421.5	24.13	414.4
19.5	24.92	401.0	25.36	394.2
20.5	26.15	382.4	26.60	376.0

of the  $^2\Pi$  ClO ground electronic state, the  $^2\Pi_{1/2}$  levels are not Zeeman active (because  $\Lambda = \pm 1$ ,  $\Sigma = \mp 1/2$ ), while the  $^2\Pi_{3/2}$  ( $\Lambda = \pm 1$ ,  $\sigma = \pm 1/2$ ) levels are Zeeman active. In addition, from eq. (10) one sees that the lower  $J$  transitions exhibit a greater range of Zeeman tuning, and the requirement for a close match with FIR laser frequency is therefore less stringent. Unfortunately the longest wavelength laser emission attainable in our spectrometer was about  $900\ \mu\text{m}$ , limiting the attempt of ClO detection to the  $J = 9.5 \leftarrow 8.5$  and higher transitions.

The LMR spectra of  $^{35}\text{ClO}$  and  $^{37}\text{ClO}$  have been observed at five different FIR laser frequencies as shown in Figures 1 and 2. Identical spectra were recorded by reacting Cl with  $\text{Cl}_2\text{O}$  or  $\text{O}_3$ , or by reacting O atoms with a large excess of  $\text{Cl}_2$ . Thus the identity of the absorbing species is certainly ClO. Table II lists all the ClO rotational transitions detected, the FIR laser wavelengths, laser gases, and the laser electric vector polarization with respect to the magnetic field ( $\pi$  = parallel,  $\sigma$  = perpendicular). The assignments were made based on the molecular constants reported previously [17].

For the  $J' = 11.5 \leftarrow J'' = 10.5$   $\pi$ -polarization transition of  $^{37}\text{ClO}$  ( $^2\Pi_{3/2}$ ,  $v = 0$ ), a detailed theoretical calculation was carried out.<sup>2</sup> The calculation involves diagonalization of an effective Hamiltonian matrix ( $\mathcal{H}_{\text{eff}}$ ) between basis-set wave functions to obtain rotational energies as a function of magnetic field strength. In this study only the  $^2\Pi_{3/2}$  wave functions are

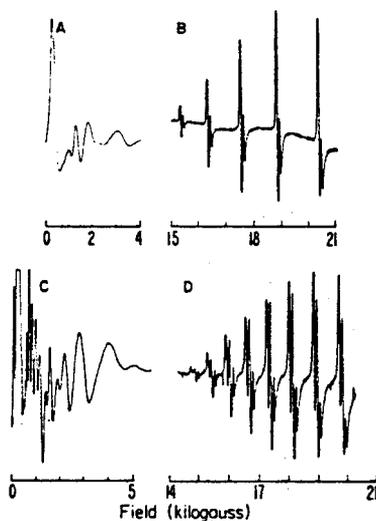


Figure 1. LMR spectra of ClO. A— $\pi$  spectrum of  $^{35}\text{ClO}$  ( $J = 14.5 \leftarrow 13.5$ ) using  $556.9\ \mu\text{m}$   $\text{CD}_3\text{I}$  line; B— $\sigma$  spectrum of  $^{35}\text{ClO}$  ( $J = 11.5 \leftarrow 10.5$ ) using  $699.5\ \mu\text{m}$   $\text{CH}_3\text{OH}$  line; C—same as A, except in  $\sigma$  polarization; D— $\sigma$  spectrum of  $^{37}\text{ClO}$  ( $J = 16.5 \leftarrow 17.5$ ) using  $469.1\ \mu\text{m}$   $\text{CH}_3\text{OH}$  line.

<sup>2</sup> For a complete derivation see [25].

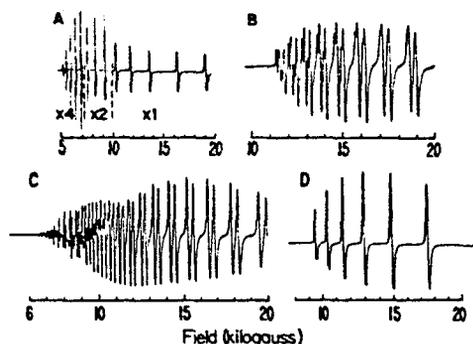


Figure 2. LMR spectra of  $^{37}\text{ClO}$ . A— $\sigma$  spectrum of  $J = 10.5 \rightarrow 11.5$  transition using  $713.1 \mu\text{m}$  DCOOH line; B— $\pi$  spectrum of  $J = 17.5 \rightarrow 18.5$  transition using  $444.4 \mu\text{m}$   $\text{CD}_3\text{I}$  line; C—same as B, except in  $\sigma$  polarization; D—same as A, except in  $\pi$  polarization.

included in the basis set because ClO is a very good Hund's case (a) type molecule, and the mixing between the  $^2\Pi_{3/2}$  and  $^2\Pi_{1/2}$  levels is not important ( $A \approx 319 \text{ cm}^{-1}$ ). The Cl nucleus has a spin of  $I = 3/2$ , and since hyperfine effects produce observable but not completely resolved splittings, nuclear spin effects are included. Thus the total angular momentum  $F = J + I$  and the basis set wave functions are characterized by the following quantum numbers:

$$(11) \quad \Psi = |\Lambda S \Sigma I m\rangle |\Phi F M_F\rangle$$

where  $S$  is the total spin angular momentum quantum number;  $m$  is the projection of  $I$  on the molecular axis and may have four values,  $3/2, 1/2,$

TABLE II. Summary of observed ClO spectra.

Species Detected	$J' + J''$	Laser Gas	$\lambda_{\text{laser}} (\mu\text{m})$	$\text{CO}_2$ Pump Laser Line	Polarization
$^{35}\text{ClO}$	11.5+10.5	$\text{CH}_3\text{OH}$	699.5	9P34	$\sigma(\Delta M_J = +1)$
$^{37}\text{ClO}$	17.5+16.5	$\text{CH}_3\text{OH}$	469.1	10R38	$\sigma(\Delta M_J = +1)$
$^{35}\text{ClO}$	14.5+13.5	$\text{CD}_3\text{I}$	556.9	10P36	$\sigma$
$^{35}\text{ClO}$	14.5+13.5	$\text{CD}_3\text{I}$	556.9	10P36	$\pi(\Delta M_J = 0)$
$^{37}\text{ClO}$	18.5+17.5	$\text{CD}_3\text{I}$	444.4	9R32	$\sigma$
$^{37}\text{ClO}$	18.5+17.5	$\text{CD}_3\text{I}$	444.4	9R32	$\pi$
$^{37}\text{ClO}$	11.5+10.5	DCOOH	713.1	10R34	$\sigma(\Delta M_J = +1)$
$^{37}\text{ClO}$	11.5+10.5	DCOOH	713.1	10R34	$\pi(\Delta M_J = 0)$

$-1/2$ , and  $-3/2$ ;  $\Phi$  is the projection quantum number of  $F$  on the molecular axis;  $\Phi = \Lambda + \Sigma + m$ ; and  $M_F$  is the projection quantum number of  $F$  in the direction of the magnetic field. The basis set consists of 32 wave functions corresponding to eight sequential values of  $F$ , and each  $F$  value generates four hyperfine sublevels corresponding to the four possible values of  $m$ .

The  $\mathcal{H}_{\text{eff}}$  is derived from a complete Hamiltonian, which is essentially the one given by Carrington [24]. Since the basis set contains only  ${}^2\Pi_{3/2}$  wave functions, this Hamiltonian is simplified to a great extent, that is, all terms that mix the  ${}^2\Pi_{3/2}$  ground state with states of different  $\Lambda$  or  $\Sigma$  are neglected. The resulting  $\mathcal{H}_{\text{eff}}$  is given below, where it has been divided into four parts, spin rotation (sr), hyperfine (hf), quadrupole (Q), and Zeeman (Z):

$$(12) \quad \mathcal{H}_{\text{sr}} = B[F(F+1) + S(S+1) + I(I+1) - \Phi^2 - \Lambda^2 - \Sigma^2 - m^2 - F_+I_- - F_-I_+] + A\Lambda\Sigma$$

where  $F_{\pm} = F_x \pm iF_y$ ,  $A$  is the spin-orbit coupling constant, and  $B$  is the rotational constant.

$$(13) \quad \mathcal{H}_{\text{hf}} = am\Lambda + (b+c)m\Sigma$$

where  $a$ ,  $b$ , and  $c$  are Frosch and Foley hyperfine constants.

$$(14) \quad \mathcal{H}_Q = \frac{e^2Qq}{4I(2I-1)} [3m^2 - I(I+1)]$$

where  $e^2Qq$  is the quadrupole coupling constant.

$$(15) \quad \mathcal{H}_Z = \mu_B H \alpha_{Zz} (\Lambda + g_e \Sigma) + g_I \mu_N H (\alpha_{Zz} + \alpha_{Zx} I_x + \alpha_{Zy} I_y)$$

where  $\alpha_{Zz}$ ,  $\alpha_{Zx}$ , and  $\alpha_{Zy}$  are direction cosine matrix elements,  $g_e$  is the electron spin  $g$  factor,  $g_I$  is the nuclear spin  $g$  factor, and  $\mu_N$  is the nuclear magneton.

The matrix elements can be calculated exactly from known relations for the direction cosine matrix elements and ladder operators (see, for example, [26]). Nonzero matrix elements of  $\mathcal{H}_{\text{eff}}$  result only for diagonal elements and for those basis set functions that are connected by the quantum number changes  $\Delta F = 0$  or  $\pm 1$  and/or  $\Delta \Phi = \Delta m = 0$  or  $\pm 1$ . A computer program which calculates these matrix elements as a function of  $F$ ,  $M_F$ , and  $H$  and diagonalizes the resulting  $32 \times 32$  matrix to yield rotational energies and wave functions has been developed [25].

The program was used to simulate the  ${}^{37}\text{ClO}$   $\pi$ -polarized,  $J = 11.5 \leftarrow 10.5$  spectrum obtained by using the  $713.1\text{-}\mu\text{m}$  DCOOH laser line. The appropriate molecular constants for  ${}^{37}\text{ClO}$  are  $B = 0.609295 \text{ cm}^{-1}$  [17],  $A = -318.7 \text{ cm}^{-1}$  [19],  $a + 1/2(b+c) = 3.1305 \times 10^{-3} \text{ cm}^{-1}$  [21],  $e^2Qq = 2.316 \times 10^{-3} \text{ cm}^{-1}$  [17] and  $g_I = 1.7395 \times 10^{-8} \text{ cm}^{-1}\cdot\text{G}^{-1}$  [27]. For this transition the energy required for zero-field absorption is  $\sim 14.009 \text{ cm}^{-1}$  (Table I). To achieve resonance with the FIR laser energy at  $\sim 14.023 \text{ cm}^{-1}$ , this rotational

transition must be Zeeman tuned by an amount equal to  $\sim 0.014 \text{ cm}^{-1}$ . The observed spectrum was simulated by calculating the magnetic fields required for resonance with each hyperfine component of a  $J', M'_J \leftarrow J'', M''_J$  transition. Although the program is written in terms of the quantum numbers  $F$  and  $M_F$ , the program output is easily interpreted in terms of  $J$  and  $M_J$  from the strong  $J$  dependence of the energies and from the relation  $M_J = M_F - M_I$ , where  $M_I$  is the projection quantum number of  $I$  on the laboratory (magnetic field) axis. The results of the calculation are given in Table III. The predicted magnetic field for each hyperfine component is also compared with the observed spectrum in Figure 3. The observed hyperfine structure is not well resolved. Only the transitions at the lower fields show structure on the high-field side of the absorption peaks. However, the observed contours are in qualitative agreement with the predicted hyperfine splittings, and the position of each hyperfine quartet is in excellent agreement with the observations. The relative intensities shown in Figure 3 are obtained by calculating a quantity proportional to the square of the dipole transition moment integral,  $\langle \Phi_u | \alpha_{Zz} |$

TABLE III. Results of calculation for the  $J' = 11.5$   
 $\leftarrow J'' = 10.5, \pi(\Delta M_J = 0)$  spectrum of  $^{37}\text{ClO}$ .

$M_J$	$H_{\text{calculated}}$ (kG)	$H_{\text{observed}}$ (kG)
-10.5	8.398	8.43
	8.370	
	8.412	
	8.510	
-9.5	9.299	9.32
	9.276	
	9.318	
	9.413	
-8.5	10.420	10.44
	10.403	
	10.442	
	10.535	
-7.5	11.855	11.87
	11.844	
	11.884	
	11.971	
-6.5	13.757	13.73
	13.752	
	13.792	
	13.875	
-5.5	16.411	16.34
	16.411	
	16.450	
	16.529	

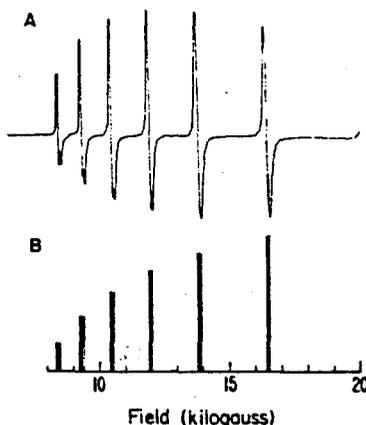


Figure 3. Comparison of the calculated results with the observed  $\pi$  spectrum of  $^{37}\text{ClO}$  ( $J = 10.5 \rightarrow 11.5$ ) using the  $713.1\text{-}\mu\text{m}$  DCOOH laser line. A—observed; B—calculated.

$\Phi_l\rangle^2$ , where the subscripts  $u$  and  $l$  refer to the upper and lower states, respectively. The relative intensity comparison between these calculated values and the LMR observations is only qualitative because other factors, such as field homogeneity and the absorption linewidth [28], affect the observed first derivative spectra. However, the overall agreement between the observed and calculated spectra clearly indicates that we are detecting ClO in the  $X^2\Pi_{3/2}(v = 0, J = 11.5 \leftarrow 10.5)$  transition.

## Kinetic Results

### *Kinetic Studies of ClO + NO<sub>2</sub> + M Reaction*

All the kinetic measurements were made under pseudo-first-order conditions ( $[\text{NO}_2] \gg [\text{ClO}]$ ) such that the  $\text{NO}_2$  concentration was constant during each measurement of the first-order ClO decay. Under such conditions, the rate expression for reaction (6) is written as

$$(16) \quad -\bar{v} \frac{d(\ln [\text{ClO}])}{dz} = k_{\text{M}}^{\text{III}}[\text{NO}_2][\text{M}] = k^{\text{II}}[\text{NO}_2] = k^{\text{I}}$$

where  $\bar{v}$  is the average flow velocity,  $z$  is the distance from the end of the movable injector to the detection volume,  $[\text{M}]$  is the concentration of the third body, and  $k^{\text{I}}$ ,  $k^{\text{II}}$ , and  $k^{\text{III}}$  are the effective first-order, second-order, and third-order rate constants, respectively.

A set of typical ClO decay plots is shown in Figure 4. Each measurement was made by varying the reaction distance  $z$  by 20–30 cm. The decay of  $[\text{ClO}]$  in the measurements was greater than a factor of 1.5, but less than

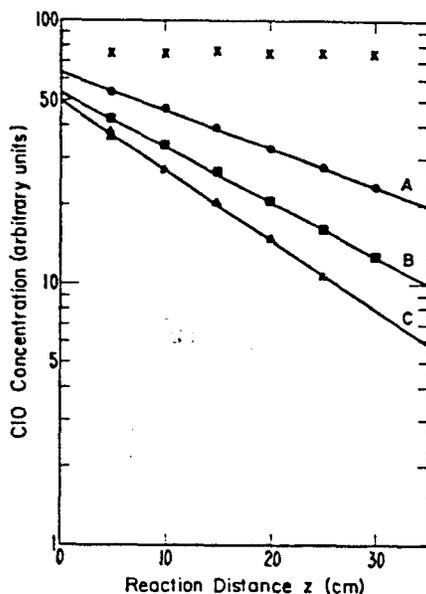


Figure 4. Pseudo-first-order decay plot for  $\text{ClO} + \text{NO}_2 + \text{O}_2$  reaction.  $T = 298$  K,  $P_t = 3.09$  torr,  $\bar{v} = 480$  cm/s,  $[\text{O}_2] = 8.46 \times 10^{16}$  molec/cm<sup>3</sup>,  $[\text{He}] = 1.44 \times 10^{16}$  molec/cm<sup>3</sup>. A— $[\text{NO}_2] = 8.04 \times 10^{14}$  molec/cm<sup>3</sup>; B— $[\text{NO}_2] = 1.14 \times 10^{15}$  molec/cm<sup>3</sup>; C— $[\text{NO}_2] = 1.50 \times 10^{15}$  molec/cm<sup>3</sup>. The X symbols are the ClO blank measurements showing negligible ClO wall loss,  $[\text{NO}_2] = 0$ .

a factor of 15. The slope of the semilogarithmic plot multiplied by  $\bar{v}$  gives the uncorrected first-order rate constant,  $k_0^I$  (in s<sup>-1</sup>). The rate constant was corrected for axial diffusion<sup>3</sup>:

$$(17) \quad k^I = k_0^I(1 + Dk_0^I/\bar{v}^2)$$

where  $D$  is the diffusion coefficient in cm<sup>2</sup>/s. The ClO diffusion coefficients were estimated as  $490(T/300)^{1.7}P^{-1}$  torr-cm<sup>2</sup>/s ( $P$  in torr) in He and  $130(T/300)^{1.7}P^{-1}$  in N<sub>2</sub> and O<sub>2</sub>. The axial diffusion correction was small, generally less than 6%.

The X symbols near the top of Figure 4 indicate the ClO concentration measured with no NO<sub>2</sub> added to the flow system. The near zero change in  $[\text{ClO}]$  at  $T > 250$  K demonstrates that the ClO loss on the flow tube surface is negligible,  $k_w < 0.1$  s<sup>-1</sup>, and that there are no other reactions of ClO.

Figure 5 shows a plot of  $k^I$  versus  $[\text{NO}_2]$  for the  $\text{ClO} + \text{NO}_2 + \text{O}_2$  reaction. The results obtained from Figure 4, lines A–C, are marked A, B, and C, respectively. At each temperature  $[\text{NO}_2]$  was varied while the rest of the experimental conditions remained unchanged. The excellent linearity and zero intercept of this plot show that the equation  $k^I = k^{II}[\text{NO}_2]$  is valid,

<sup>3</sup> The axial diffusion correction is given in [29a]; the diffusion coefficients are estimates based on data given in [29b].

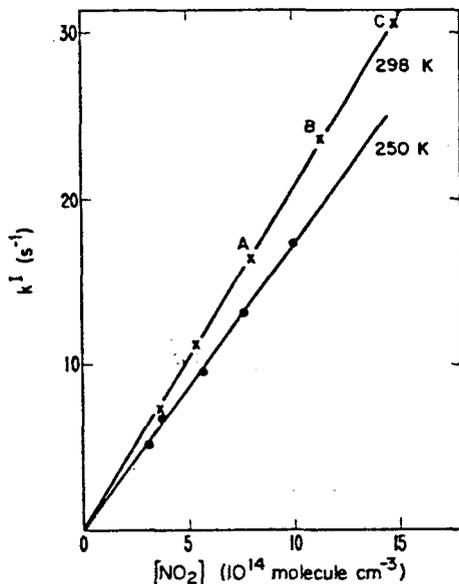


Figure 5. Plot of  $k^I$  versus  $[\text{NO}_2]$  for  $\text{ClO} + \text{NO}_2 + \text{O}_2$  reaction.  $\times$ — $T = 298 \text{ K}$ ,  $P_t = 3.09 \text{ torr}$ , slope =  $2.04 \times 10^{-14} \text{ cm}^3/\text{molec}\cdot\text{s}$ ;  $\bullet$ — $T = 250 \text{ K}$ ,  $P_t = 1.24 \text{ torr}$ , slope =  $1.71 \times 10^{-14} \text{ cm}^3/\text{molec}\cdot\text{s}$ .  $[\text{NO}_2]$  at this temperature has been corrected for dimerization. The results from Figure 4 are marked A, B, and C, respectively.

and the slope gives  $k^{II}$ . Some  $k^{II}$  measurements in this study were obtained by simply dividing the measured  $k^I$  value by  $[\text{NO}_2]$ .

A plot of  $k^{II}$  versus  $[M]$  (total gas density) at 250 and 298 K using He or  $\text{O}_2$  as a carrier gas is shown in Figure 6. The data shown in Figure 5 for  $T = 298$  and 250 K are indicated by A and B. The linearity of this plot demonstrates that the reaction is third order at pressures below 6 torr. For

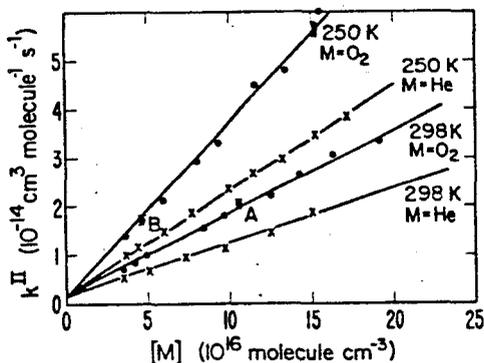


Figure 6. Plot of  $k^{II}$  versus  $[M]$  for the  $\text{ClO} + \text{NO}_2 + M$  reactions at 250 and 298 K.  $\bullet$ — $M = \text{O}_2$ ;  $\times$ — $M = \text{He}$ . The data shown in Figure 5 are indicated by A and B.

experiments using He as the third body, the slope gives  $k_{\text{He}}^{\text{III}}$  directly, because the concentrations of other species such as  $\text{NO}_2$ ,  $\text{O}_2$ , and  $\text{Cl}_2$  is negligible. However, for experiments using  $\text{N}_2$  or  $\text{O}_2$  as the third body the amount of He carrier added through the microwave discharge (ClO source) and the ozone reservoir comprises up to 30% of the gas mixture in the flow tube at low pressures. To derive accurate  $k_M^{\text{III}}$  values for various third bodies, a linear least-square fit of the data to the expression

$$(18) \quad k^{\text{II}} = k_w^{\text{II}} + k_{\text{He}}^{\text{III}}[\text{He}] + k_{\text{O}_2}^{\text{III}}[\text{O}_2] + k_{\text{N}_2}^{\text{III}}[\text{N}_2]$$

was carried out by computer. The intercept in the  $k^{\text{II}}$  versus  $[M]$  plot,  $k_w^{\text{II}}$ , is a pressure-independent heterogeneous contribution to  $k^{\text{II}}$ , which can be taken as the second-order rate constant for the  $\text{ClO} + \text{NO}_2$  reaction with the stabilization by wall collisions. Thus for a simultaneous fit, the  $k_w^{\text{II}}$  value is constrained to be the same for different carrier gases at the same temperature. The magnitude of  $k_w^{\text{II}}$  is relatively small,  $(0.5\text{--}2.1) \times 10^{-15} \text{ cm}^3/\text{molec}\cdot\text{s}$ , and varies with different wall conditions and temperatures. Therefore experimental data obtained at different times with different wall conditions were grouped separately, as shown in Table IV.

In Table IV, the experimental conditions,  $k_w^{\text{II}}$ , and  $k_M^{\text{III}}$  values for  $\text{N}_2$  at 297 K and for He and  $\text{O}_2$  at six temperatures ranging from 250 to 416 K are summarized. The listed uncertainties for  $k_w^{\text{II}}$  and  $k_M^{\text{III}}$  are one standard deviation of the coefficients derived from the fits of the experimental data to eq. (18). At room temperature  $k_{\text{He}}^{\text{III}}$  was measured as  $(1.02 \pm 0.02)$ ,  $(1.06 \pm 0.02)$ , and  $(1.07 \pm 0.02) \times 10^{-31} \text{ cm}^6/\text{molec}^2\cdot\text{s}$  at three different times with

TABLE IV. Data summary for  $\text{ClO} + \text{NO}_2 + M$  reaction.

T (°K)	M	No. Exp.	$P_{\text{He}}$ (torr)	$P_{\text{N}_2}$ (torr)	$[\text{NO}_2]$ ( $10^{15} \text{ molec}/\text{cm}^3$ )	$k_w^{\text{II}}$ ( $10^{-31} \text{ cm}^6/\text{molec}^2\cdot\text{s}^{-1}$ )	$k_M^{\text{III}}$ ( $10^{-15} \text{ cm}^3/\text{molec}^2\cdot\text{s}^{-1}$ )
296	He	37	0.91-6.60	0.91-6.60	0.50-3.75	1.02±0.02	2.1±0.3
	$\text{N}_2$	14	0.67-3.70	0.12-0.48	0.94-1.86	2.06±0.06	
296	He	6	1.10-5.92	1.10-5.92	0.37-2.08	1.06±0.02	1.5±0.2
	$\text{O}_2$	11	0.56-5.57	0.26-0.47	0.43-1.68	1.85±0.02	
296	He	18	0.86-4.88	0.86-4.88	0.65-2.34	1.07±0.03	1.9±0.2
	$\text{O}_2$	14	0.96-4.17	0.25-0.38	0.53-1.81	1.88±0.03	
	$\text{N}_2$	8	0.87-4.71	0.24-0.33	0.40-1.23	2.12±0.03	
250	He	9	0.97-4.46	0.97-4.46	0.38-0.83	2.26±0.07 <sup>a</sup>	0.7±0.1
	$\text{O}_2$	17	0.71-4.53	0.21-0.29	0.19-1.04	3.86±0.06 <sup>a</sup>	
272	He	8	1.13-4.47	1.13-4.47	0.26-0.80	1.50±0.06	1.6±0.6
	$\text{O}_2$	8	0.72-4.23	0.19-0.23	0.25-0.67	2.73±0.06	
342	He	9	0.95-4.46	0.95-4.46	0.80-1.72	0.72±0.02	0.5±0.1
	$\text{O}_2$	10	0.57-4.16	0.21-0.34	0.92-1.69	1.16±0.02	
387	He	8	1.03-3.97	1.03-3.97	1.18-1.82	0.46±0.03	0.9±0.2
	$\text{O}_2$	10	0.24-0.29	0.24-0.29	1.48-1.76	0.74±0.03	
416	$\text{O}_2$	7	0.71-3.58	0.31-0.40	1.33-1.85	0.58±0.03 <sup>b</sup>	0.3±0.1

<sup>a</sup> Measured with correction for  $\text{NO}_2$  dimerization. Without correction,  $k_{\text{He}}^{\text{III}}$  and  $k_{\text{O}_2}^{\text{III}}$  were 2.17 and  $3.69 \times 10^{-31} \text{ cm}^6/\text{molec}^2\cdot\text{s}$ , respectively.<sup>2</sup> See text.

<sup>b</sup> Assuming  $k_{\text{He}}^{\text{III}} = 3.8 \times 10^{-32} \text{ cm}^6/\text{molec}^2\cdot\text{s}$ . See text.

excellent consistency. Thus an average value of  $1.05 \times 10^{-31} \text{ cm}^6/\text{molec}^2\text{-s}$  is obtained for the  $\text{ClO} + \text{NO}_2 + \text{He}$  reaction at room temperature. Similarly the room temperature values of  $1.87 \times 10^{-31}$  and  $2.09 \times 10^{-31} \text{ cm}^6/\text{molec}^2\text{-s}$  are obtained for  $M = \text{O}_2$  and  $M = \text{N}_2$ , respectively.

At 416 K the  $\text{ClO} + \text{NO}_2 + \text{He}$  reaction was too slow to be measured. The data for  $M = \text{O}_2$  at this temperature were corrected for the small contribution from  $M = \text{He}$  using  $k_{\text{He}}^{\text{III}} = 3.8 \times 10^{-32} \text{ cm}^6/\text{molec}^2\text{-s}$  obtained by extrapolating the lower temperature data.

In Figure 7 the data for  $k_{\text{He}}^{\text{III}}$  and  $k_{\text{O}_2}^{\text{III}}$  are plotted as a function of  $1/T$ . The least-square fits gave  $k_{\text{He}}^{\text{III}} = (2.79 \pm 0.38) \times 10^{-33} \exp[(1087 \pm 41)/T]$  and  $k_{\text{O}_2}^{\text{III}} = (3.48 \pm 0.22) \times 10^{-33} \exp[(1183 \pm 19)/T] \text{ cm}^6/\text{molec}^2\text{-s}$ , where the error limits are one standard deviation limits on the coefficients.

As mentioned earlier, the concentration of  $\text{N}_2\text{O}_4$  becomes significant only at the lowest temperatures. An attempt was made to measure the rate constant for the reaction of  $\text{ClO}$  with  $\text{N}_2\text{O}_4$ . A special double port movable inlet was constructed. One port was designed to maximize the  $\text{N}_2\text{O}_4$  concentration injected into the flow tube. It had a short ( $\approx 1.5 \text{ cm}$ ) length of 0.09-mm inside diameter capillary attached to the end of a 14-mm inside diameter  $\times$  10-cm length bulb. This constriction caused a high pressure ( $\sim 40 \text{ torr}$ ) to develop inside the inlet and therefore enhanced the  $\text{NO}_2$  dimerization. The pressure in the inlet was measured and the fractional dimerization was calculated using the equilibrium constant [15]. At 255

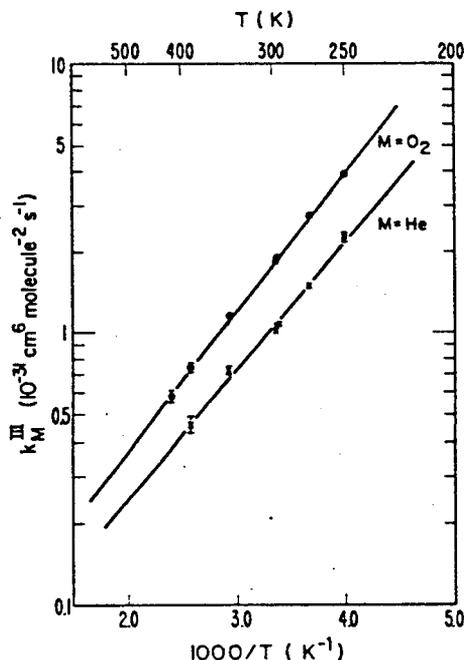


Figure 7. Arrhenius plot of  $k_{\text{O}_2}^{\text{III}}$  and  $k_{\text{He}}^{\text{III}}$  for  $\text{ClO} + \text{NO}_2 + M$  reaction.

K the  $N_2O_4$  flow rate was typically 70–80% of the total flow rate ( $N_2O_4 + NO_2$ ). Equilibrium was assured by the long residence time in the cooled inlet tube ( $\sim 7$  s). The second port was designed for minimum dimerization, with low pressure and a short residence time. For these measurements the ClO was added to the carrier gas upstream from the reaction zone. When a mixture of about 75%  $N_2O_4$  and 25%  $NO_2$  was reacted with ClO, the observed ClO loss rate was significantly less ( $\sim 1/2$ ) than when the same flow rate was added through the second injector port as monomers. From this observation we conclude that small amounts of  $N_2O_4$  do not interfere with the ClO +  $NO_2$  kinetic measurements. The  $N_2O_4$  is probably not dissociated significantly in the flow tube at 255 K and pressures of a few torr. If that is true, then  $k(\text{ClO} + N_2O_4) \approx 2 \times 10^{-14} \text{ cm}^3/\text{molec}\cdot\text{s}$ .

### Kinetic Studies of $\text{ClO} + \text{NO} \rightarrow \text{Cl} + \text{NO}_2$ Reaction

A typical ClO decay plot is shown in Figure 8. In these measurements the reaction distance  $z$  was varied by 15–30 cm. The ClO decay with no NO added ( $\times$  symbols at the top of Figure 8) also indicated negligible wall loss. The axial diffusion correction was applied and was less than 8% at 209 K and 3% at 393 K.

Figure 9 shows a plot of  $k^1$  as a function of  $[\text{NO}]$  for data at 393, 296, and 202 K. At room temperature, some measurements (indicated  $\blacktriangle$  on Figure

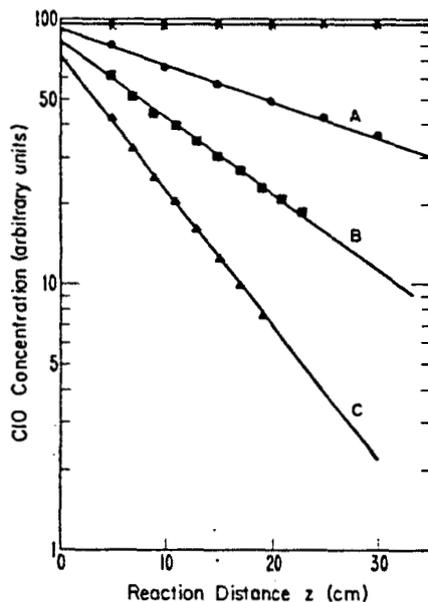


Figure 8. Pseudo-first-order decay plot for ClO + NO reaction.  $T = 393 \text{ K}$ ,  $P_t = 1.91 \text{ torr}$ ,  $\bar{v} = 2210 \text{ cm/s}$ . A— $[\text{NO}] = 5.27 \times 10^{12} \text{ molec/cm}^3$ ; B— $[\text{NO}] = 1.19 \times 10^{13} \text{ molec/cm}^3$ ; C— $[\text{NO}] = 2.01 \times 10^{13} \text{ molec/cm}^3$ . The  $\times$  symbols are the ClO measurements showing negligible ClO wall loss,  $[\text{NO}] = 0$ .

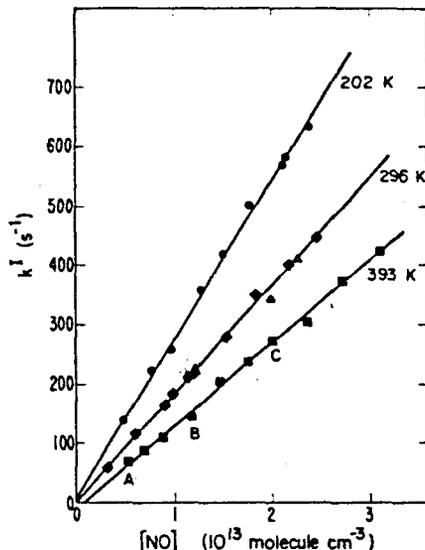


Figure 9. Plot of  $k^I$  versus  $[NO]$  at 202, 296, and 393 K for  $ClO + NO$  reaction. At 296 K.  $\blacklozenge$ —NO through the movable injector;  $\blacktriangle$ —ClO through the movable injector. A, B, and C indicate data from Figure 8.

9) were performed by adding ClO through the movable inlet and adding NO with the carrier gas. No difference in rate constant was observed. No pressure dependence was observed over the range of 1–3 torr. The data for  $k^I$  versus  $[NO]$  at each temperature were fitted with a least-squares program to a straight line. The slope  $k^{II}$ , intercept, and experimental conditions at five temperatures ranging from 202 to 393 K are summarized in Table V. At room temperature the rate constants were measured as 1.66, 1.84, and  $1.84 \times 10^{-11} \text{ cm}^3/\text{molec}\cdot\text{s}$  at three different times. These give an average value of  $1.78 \times 10^{-11} \text{ cm}^3/\text{molec}\cdot\text{s}$ .

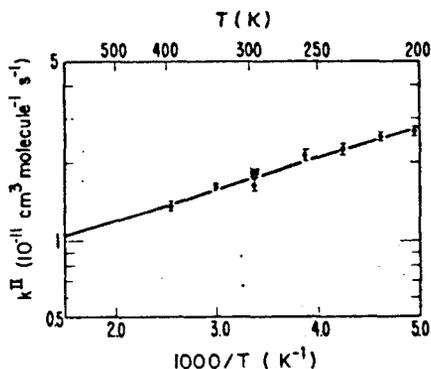


Figure 10. Arrhenius plot for  $ClO + NO \rightarrow Cl + NO_2$  reaction rate constant data.

TABLE V. Data summary for ClO + NO reaction.

T (°K)	No. Exp	$P_r$ (torr)	[NO] ( $10^{13}$ molecule $\text{cm}^{-3}$ )	$k$ ( $10^{-11}$ $\text{cm}^3$ molecule $^{-1}$ $\text{s}^{-1}$ )	Intercept ( $\text{s}^{-1}$ )
295	10	0.98-2.56	0.33-2.47	1.84±0.03	0±4
296	17	1.05-3.10	0.82-2.97	1.66±0.04	13±3
297	8	1.11	0.41-2.10	1.84±0.03	10±5
202	10	1.06-1.99	0.49-2.39	2.64±0.06	14±9
217	9	1.01-1.94	0.33-2.40	2.52±0.05	13±7
236	9	1.04-1.95	0.43-2.30	2.25±0.06	7±8
258	10	1.11-1.12	0.92-3.19	2.13±0.05	-13±12
333	10	1.01-2.61	0.36-2.77	1.63±0.02	-1±4
393	10	1.91-2.77	0.53-3.10	1.39±0.03	-11±6

The measured rate constants are plotted versus  $1/T$  in Figure 10. A least-squares fit to the Arrhenius expression gives  $k_3 = (7.13 \pm 0.48) \times 10^{-12} \exp [(271 \pm 18)/T] \text{ cm}^3/\text{molec}\cdot\text{s}$ , where the error limits represent one standard deviation.

### Discussion

The measurements of  $k(\text{ClO} + \text{NO}_2 + \text{M})$  are summarized with the work from other laboratories in Table VI. The error limits have been increased to allow for systematic errors. Our  $\text{M} = \text{He}$  measurement is about in the middle of the values obtained by Leu et al. [8] and Zahniser et al. [10] and agrees well with the 297 K value reported by Birks et al. [9]. The range of these measurements at room temperature is about  $\pm 15\%$ . For  $\text{M} = \text{N}_2$  we measured  $k$  only at 297 K. Our result is higher than all of the published values. The maximum discrepancy is 38% compared to Zahniser et al. Although the error bars on these measurements nearly overlap, this dif-

TABLE VI. Comparison of results for ClO + NO<sub>2</sub> + M reaction.

M	Arrhenius Form	T (K)	$k(298\text{K})$ ( $10^{-11}$ $\text{cm}^3$ molecule $^{-1}$ $\text{s}^{-1}$ )	P (torr)	Technique <sup>a</sup>	Reference
He	$(2.8 \pm 0.6) \times 10^{-33} \exp[(1090 \pm 80)/T]$	250-387	1.05±0.15	0.67-6.60	DF-LMR	This Work
	$(3.54 \pm 0.06) \times 10^{-33} \exp(950/T)$	250-365	0.83±0.12	1.1-6.5	DF-RF	Zahniser et al. <sup>10</sup>
	$(2.66 \pm 0.35) \times 10^{-33} \exp[(11140 \pm 40)/T]$	248-417	1.20±0.05	1-9	DF-MS	Leu et al. <sup>8</sup>
	-	297	0.96±0.03	2.5-5.0	DF-MS	Birks et al. <sup>9</sup>
N <sub>2</sub>	-	297	2.09±0.30	0.67-3.70	DF-LMR	This work
	-	297	1.52±0.23	2-4.1	DF-RF	Zahniser et al. <sup>10</sup>
	$(3.69 \pm 0.24) \times 10^{-33} \exp[(1150 \pm 20)/T]$	298-417	1.72±0.17	1-6	DF-MS	Leu et al. <sup>8</sup>
	$(4.40 \pm 0.86) \times 10^{-33} \exp[(1087 \pm 70)/T]$	250-356	1.85±0.05	1.5-4.5	DF-MS	Birks et al. <sup>9</sup>
	$(4.3 \pm 1.2) \times 10^{-33} \exp[(1085 \pm 66)/T]$	274-339	1.64	50	MA	Cox and Lewis <sup>30</sup>
		1.5 ± 0.12	20	FP-A	Molina et al. <sup>34</sup>	
O <sub>2</sub>	$(3.5 \pm 0.6) \times 10^{-33} \exp[(1180 \pm 80)/T]$	250-416	1.86±0.28	0.56-5.57	DF-LMR	This work
Ar	-	298	1.15±0.10	1-4	DF-MS	Leu et al. <sup>8</sup>

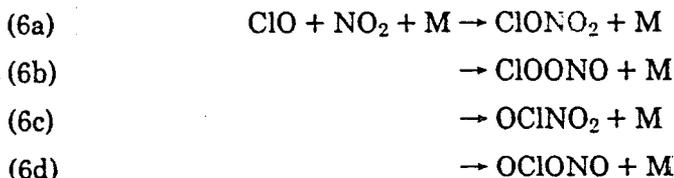
<sup>a</sup> DF—discharge flow; LMR—laser magnetic resonance; RF—resonance fluorescence; MS—mass spectrometry; MA—modulated absorption; FP-A—flash photolysis ultraviolet absorption.

ference is larger than is normally found. This work is the only reported measurement with  $M = O_2$ . We can compare this result with the earlier work on the basis of the temperature dependence, which agrees very well with previous measurements in He and  $N_2$ .

Cox and Lewis [30] have reported the most extensive pressure dependence study of reaction (6). They measured the rate constant in  $N_2$  over the range of 25–612 torr using modulated photolysis with ultraviolet absorption detection of ClO. Their data connect reasonably well with the low-pressure data, as seen in Table VI. They report some falloff at pressures greater than 50 torr.

Molina et al. [34] have measured  $k_6$  at 298 K in 20 and 100 torr  $N_2$  using a flash-photolysis ultraviolet absorption technique. Their result using a  $Cl_2O$  source of ClO agrees reasonably well with the others.

The product of the termolecular reaction of ClO and  $NO_2$  was not quantitatively observed here or in most of the previous studies [8–10]. It has been suggested [31–34] that more than one isomer of  $ClONO_2$  may be formed. The possibilities include the following



The studies of the thermal decomposition of  $ClONO_2$  [31] and the isotopic exchange reaction between  $ClONO_2$  and  $^{15}NO_2$  by Knauth and co-workers [32], when combined with the equilibrium constant for  $ClONO_2$  formation, lead to a value of  $k_{N_2}^{III}$  that is about 30% of the values listed in Table VI. It is not clear that this difference is significant due to uncertainties in the thermochemical data. Chang et al. [33] suggested, on the basis of a recombination theory calculation, that if the isomer  $ClOONO$  is formed along with  $ClONO_2$ , the discrepancy between the forward and reverse rate constants of  $ClONO_2$  can be resolved. Recently, in a flash-photolysis experiment, Molina et al. [34] have observed different rates of removal of ClO depending upon the ClO precursor ( $Cl_2O$  or  $OCIO$ ). They proposed that an  $OCIONO$  isomer is formed and takes part in a secondary reaction. Using infrared absorption spectroscopy, they also found that the amount of  $ClONO_2$  produced was only about one third of the amount of ClO removed. Thus it is probable that more than one product is formed in this reaction. This conclusion has a very important implication to atmospheric chemistry because the identity of the product molecule determines the photolysis rates, and products and hence the effect of the  $ClO + NO_2 + M$  reaction. The large rate constant for the reaction makes it a potentially important process for ClO in the stratosphere. Therefore the reaction products must be determined and evaluated.

The measurement using O<sub>2</sub> as a carrier gas was carried out to investigate the possibility of a reaction rate enhancement by O<sub>2</sub>. The reaction of O<sub>2</sub> with ClO to form OClO<sub>2</sub>,



may be possible,<sup>4</sup> although no information on the reaction rate is available. If reaction (7) is sufficiently fast and followed by the reaction



the rate of ClNO<sub>3</sub> formation could be enhanced in oxygen. Assuming a steady state of [ClO<sub>2</sub>], the rate constant enhancement is  $k_{19}k_7[\text{O}_2]/(k_{-7}[\text{M}] + k_{19}[\text{NO}_2])$ , where  $k_{-7}$  is the back-reaction rate constant for reaction (7). In the present experiment the variation of [O<sub>2</sub>] and [NO<sub>2</sub>] was not large enough to observe a significant enhancement due to this mechanism.

The measured relative third-body efficiencies of 1.0:1.8:2.0 for He, O<sub>2</sub>, and N<sub>2</sub>, respectively, in ClO + NO<sub>2</sub> + M reaction can be compared with those obtained for the similar reaction



by Howard [36], namely, 1.0:1.5:2.1. Although the O<sub>2</sub> third-body efficiency measured in this study for reaction (6) is slightly higher than that of reaction (20), the difference is too small to be considered significant.

The activation energy of  $k_{\text{O}_1}^{\text{III}}$ , (1180 ± 80), is slightly higher than that of  $k_{\text{He}}^{\text{III}}$ , (1090 ± 80). Although the difference may be due to the experimental uncertainties, the possibility remains that this is an indication of a different mechanism for ClO + NO<sub>2</sub> reaction in O<sub>2</sub>. We plan to extend these measurements to higher pressures as a further test for a mechanism involving OClO<sub>2</sub>.

The measurements of the ClO + NO rate constant  $k_3$  are summarized

TABLE VII. Comparison of rate constant measurements for ClO + NO → Cl + NO<sub>2</sub> reaction.

Temp. Range (K)	Arrhenius Form (cm <sup>3</sup> molecule <sup>-1</sup> s <sup>-1</sup> )	k(298K) (10 <sup>-11</sup> cm <sup>3</sup> molecule <sup>-1</sup> s <sup>-1</sup> )	Method <sup>a</sup>	Reference
202-393	$(7.1 \pm 1.4) \times 10^{-12} \exp[(270 \pm 50)/T]$	1.78 ± 0.27	DF-LMR	This work
298	-	1.7 ± 0.2	DF-MS	Clyne and Watson <sup>3</sup>
230-295	$(1.13 \pm 0.14) \times 10^{-11} \exp[(200 \pm 30)/T]^b$	2.2 ± 0.4 <sup>b</sup>	DF-RF	Zahniser and Kaufman <sup>4</sup>
227-415	$(5.72 \pm 0.18) \times 10^{-12} \exp[(296 \pm 20)/T]$	1.53 ± 0.11	DF-MS	Lau and DeMore <sup>5</sup>
298	-	1.58 ± 0.16	DF-MS	Clyne and MacRobert <sup>6</sup>
298	-	1.75 ± 0.15	DF-MS	Ray and Watson <sup>7</sup>

<sup>a</sup> DF—discharge flow; LMR—laser magnetic resonance; MS—mass spectrometry; RF—resonance fluorescence.

<sup>b</sup> Based on a rate constant ratio measurement as described in text.

<sup>4</sup> Prasad [35] also recently discussed some possible implications of OClO<sub>2</sub> in stratospheric chemistry.

in Table VII. Our room temperature measurement is in excellent agreement with those of Clyne and Watson [3] and Ray and Watson [7] and in good agreement with the other measurements. The temperature dependence is very close to the result of Leu and DeMore [5], although their rate constants fall about 15% below ours. The temperature dependence given by Zahniser and Kaufman [4] is also in good agreement, but their values fall somewhat higher than ours. The  $k_3$  given in Table VII for Zahniser and Kaufman is based on their measurement of the ratio of  $k_3$  to  $k(\text{Cl} + \text{O}_3)$  and their value for  $k(\text{Cl} + \text{O}_3)$ .

In conclusion, we find that the LMR technique can be used to detect ClO radicals with sufficient sensitivity to allow pseudo-first-order kinetic studies in a flow-tube reactor. This method offers the advantages of positive spectroscopic identification of the ClO with no interference from product species or from problems encountered with indirect measurement methods.

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