

Characterization of a Cold Cesium Source for PARCS: Primary Atomic Reference Clock in Space

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Abstract

The PARCS (Primary Atomic Reference Clock in Space) project is a joint NIST-JPL-University of Colorado venture aimed at placing a cesium (Cs) atomic clock aboard the International Space Station (ISS). This orbiting clock will achieve high accuracy, in part due to the long Ramsey times afforded by the microgravity environment, and allow for precision tests of fundamental physics including relativity theory. As part of this effort, we are evaluating the characteristics of a prototype cold Cs source based on launching atoms from an optical molasses. The apparatus described here will be used to develop other PARCS components such as the microwave cavity structure and detection systems, and to investigate two-dimensional cooling schemes for future Cs fountains and space clocks.

Introduction

Laser cooling techniques have allowed the development of modern fountain clocks that now achieve uncertainties of $\approx 1 \times 10^{-15}$. The maximum practical Ramsey time achievable in a fountain is on the order of 1 s. Recently, several groups have begun the development of space based clocks in order to obtain the longer Ramsey times (~ 10 s) possible in a microgravity environment, to perform tests of fundamental physics such as relativity theory, and to provide a globally available realization of the second.

Our group at NIST is collaborating with JPL and the University of Colorado to develop PARCS (Primary Atomic Reference Clock in Space), a cesium (Cs) atomic clock that is scheduled to fly aboard the International Space Station (ISS).

An overview of the PARCS space clock project has been presented elsewhere [1]. Similarly, the French PHARAO

space clock addresses many of the same issues and problems [2]. Our present goal is to achieve a stability of $\sigma_y(\tau) \approx 5 \times 10^{-14} \tau^{-1/2}$ with an ultimate uncertainty of $\sim 10^{-16}$. The final clock parameters will depend on many factors, some of which are still under study. PARCS will operate in the configuration illustrated in Fig. 1.

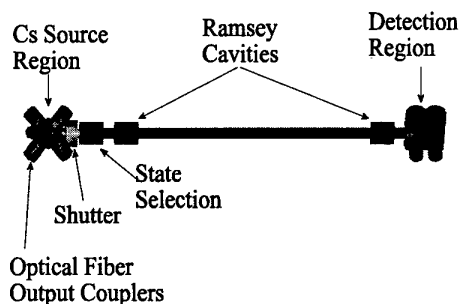


Figure 1. Proposed configuration for PARCS.

Cold Cs atoms are launched from an optical molasses, state selected, then enter a Ramsey-type microwave cavity. The shutters allow the clock to operate in a multiple launch mode; that is, a second ball of atoms can be loading in the source region while the previous ball is in the interrogation region, thus increasing the flux of atoms through the clock. Upon exiting the interrogation region, the atoms enter a detection region where the populations in $F=3$ and $F=4$ are measured.

We have developed theoretical models of atom flux requirements that yield a range of operating parameters which

achieve $\sigma_y(\tau) \approx 5 \times 10^{-14} \tau^{-1/2}$. Consider a clock with Ramsey cavities separated by 75 cm. With a molasses load time of 0.3 s, an atom temperature of 2 μ K, a Ramsey time of 5 s, and 4 launches per lineside, we require $\approx 2 \times 10^6$ atoms in the $m_f=0$ clock state per launch. The same conditions with a 10 s Ramsey time boosts required atom number to 5×10^6 per launch.

Experimental Apparatus

Our experimental apparatus is shown in Fig. 2. The vacuum chamber consists of the source region, a long toss tube (~ 1.5 m), and a detection region. The source region contains twelve 3.8 cm windows. Six windows are located on the faces of a cube oriented in the (1,1,1) beam geometry [3]. These are used for introducing the cooling beams into the vacuum. The remaining six windows are for viewing and for future experimental studies. There are several smaller ports containing a Cs oven and a small window suitable for introducing a Zeeman slower beam. This system contains three vacuum pumps an ion pump located at the top of the toss tube, graphite getters located near the center of the source, and a Ti sublimation pump between the source and the toss tube. The source region is surrounded by three pairs of Helmholtz coils for zeroing the ambient magnetic field.

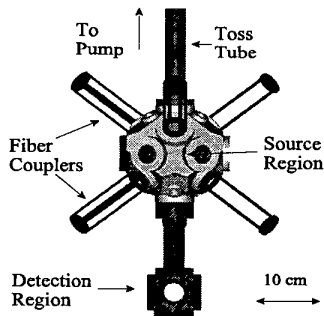


Figure 2. The cold Cs source experimental apparatus.

An optical amplifier pumped by a DBR (Distributed Bragg Reflector) diode laser supplies ~ 200 mW of light at 852 nm. The laser light is passed through an optical isolator and then divided into two beams that are delivered to separate double-pass acoustic-optic modulators (AOM) allowing the three upward pointing beams to be shifted in frequency independently from the three downward beams for launching. Nominally, the beams are shifted by the AOMs and are detuned $\sim 2\Gamma$ red of the $F=4 \rightarrow F'=5$ transition. The beams are then further split and approximately equal powers are coupled into 6 polarization-maintaining fiber-optic cables. The fibers extend to output couplers mounted on the chamber windows that expand the beams to a $1/e^2$ diameter of ~ 1.5 cm. The fiber output couplers contain polarizing beamsplitting cubes

that direct the beam originating from the opposite output coupler onto a photodiode. This allows the intensity of the six cooling beams to be monitored in real time. We have developed an intensity servo for the six beams; the servo monitors the light levels at the photodiodes and then sends control signals back onto variable retarders on the optical table which in turn adjust the amount of light entering the six fiber input couplers.

An extended-cavity diode laser (ECDL) locked on the $F=3 \rightarrow F'=4$ transition serves as a repump (~ 5 mW) and this is directed onto the molasses region through one of the additional side windows. A beam from a second ECDL locked on the $F=4 \rightarrow F'=5$ transition is delivered via polarization-maintaining fiber to the vacuum system where cylindrical lenses expand the beam into a sheet of light ~ 0.5 mm thick, after which it enters a window on the detection cube that is located ~ 25 cm below the source. A retro mirror directs the detection beam back onto itself and creates a standing wave. The detection system consists of collection optics, a large photodiode and a high-gain amplifier and is located on one of the remaining windows of the detection cube. Calibration of the detection system allows the number of atoms in the ball to be determined. Additional information on atom number is provided by a CCD imaging system focused onto the molasses region.

We operate the apparatus in two basic modes, dropping atoms and launching atoms. In the drop mode, the molasses beams are turned on for a time T_{load} , the atoms are then subjected to a post-cool routine by detuning the beams to $\approx -8\Gamma$ from the $F=4 \rightarrow F'=5$ transition, and finally the beams are shut off allowing the atoms to fall to the detection region. In the launch mode, the molasses beams are on for T_{load} , the up beams are then detuned with respect to the down beams creating a moving molasses at the launch velocity. Approximately 2 msec later the atoms are post-cooled before leaving the source region. The atoms travel up, reach apogee, and fall back into the detection region.

Experimental Results

We have measured the molasses loading rate as a function of a variety of parameters. Figure 3 shows data obtained by dropping atoms from the molasses after various load times and for different beam intensities at a constant beam diameter of ~ 1.5 cm ($1/e^2$). For example, at a load time of 1 s, with intensities of ~ 5 mW/cm² per beam, $\sim 6 \times 10^7$ atoms arrive in the detection region. With the present experimental configuration, we are limited to intensities of ~ 5 mW/cm².

By recording the size of the atomic fluorescence signal in the detection region as a function of various launch heights (≥ 1 m) we can determine the temperature of the atom sample. These measurements indicate an atom temperature of ~ 12 μ K and an initial radius of ~ 0.5 cm. We suspect these relatively high atom temperatures are the result of an ineffective post-cool routine produced by fiber-coupling efficiency changes with laser detuning resulting in large beam imbalances. A

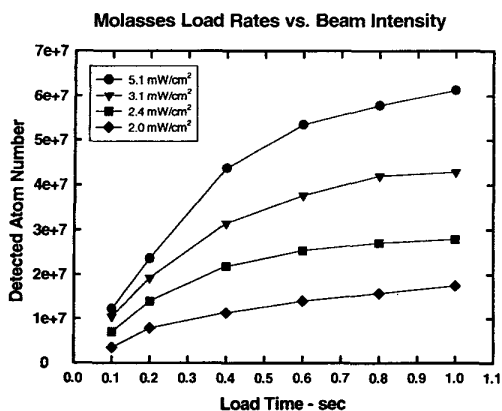


Figure 3. The number of Cs atoms arriving at the detection region after dropping from the optical molasses at various load times and beam intensities.

consequence of this high atom temperature is that there is appreciable thermal expansion of the atom ball during the interval when a sample is dropped from the source until it reaches the detection region. We estimate that approximately 1/3 of our atoms are lost at a ~ 1.3 cm diameter aperture located between the source and the detection region. This loss factor applied to the results from the dropped signal in Fig. 3 with 5 mW/cm² per beam at long load times indicates that the number originally in the source is $\sim 9 \times 10^7$. This value is in good agreement with the number of atoms detected within the molasses from fluorescence using the CCD imaging system.

We can compare the results of these measurements with the requirements outlined by our flux modeling. For example, at a load time of 0.3 s (this is the load time our present flux models incorporate) with 5 mW/cm² per beam there are $\sim 5 \times 10^7$ atoms originally in the source. Assuming equal populations among the nine m_F sub levels, this leaves $\sim 6 \times 10^6$ atoms in the $F=4$, $m_F=0$ state. According to our flux requirements model, this allows for a 10 s Ramsey times and a stability of $\sigma_y(\tau) \approx 5 \times 10^{-14} \tau^{-1/2}$ (assuming a target atom temperature of 2 μ K).

We have also characterized our Cs oven to determine how the atoms load into the molasses. In one experiment, we turned on the Cs oven and monitored the number of atoms in a dropped signal with a constant load time as the oven heated up. The oven quickly reached its set-point temperature, and during this initial warm-up the number of atoms in the signal increased at the same rate as the oven temperature. However, even after the temperature of the oven stabilized, the number of atoms in the signal continued to increase significantly for several hours. This behavior indicates that Cs atoms are loading into the molasses from both the Cs beam from the oven as well as from background vapor. Such issues are

important for the design of PARCS in which low source-region pressure is desirable to reduce collisional loss of atoms with slow launch velocities.

Conclusion

We anticipate that we will achieve colder atom samples (2 μ K) as soon as we improve control of the post-cooling process. Given these temperatures and our current atom numbers, we already meet the flux requirements to reach our target stability at ~ 10 s Ramsey times. We are continuing to investigate source operating parameters and anticipate still larger atom flux numbers can be readily achieved.

Acknowledgments

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